

## Supporting Information

### Role of hyper-reduced states in hydrogen evolution reaction at sulfur vacancy in MoS<sub>2</sub>

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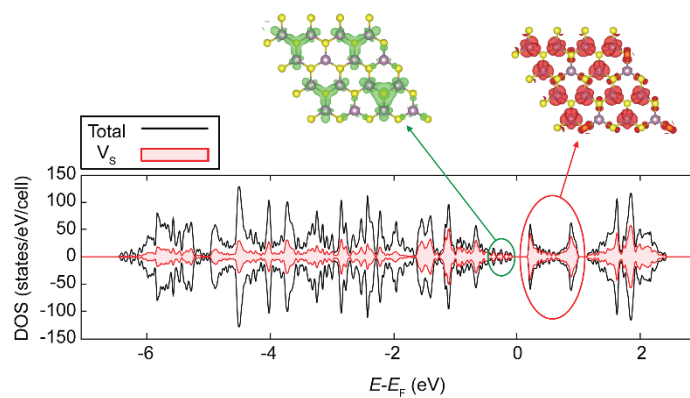
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## Detailed electronic and structural character of $V_s$ and its concentration dependence

We first explain the structural distortion and electronic properties at low vacancy densities. Once a sulfur vacancy is created, neighboring ions around the vacancy site are relaxed. The largest relaxation occurs with the three nearest Mo ions that displace inward by 0.04 Å at 2.8% of the vacancy concentration. The vacancy defect in this concentration can be regarded as an isolated one because further decrease of the concentration does not influence the defect property. S-vacancy affects electronic property of MoS<sub>2</sub> by creating defect states inside the band gap which primarily derive from *d* states of neighboring Mo ions. There are three defect states labeled as a singlet  $a_1$  and doubly degenerate  $e$  states based on the group theory. The  $a_1$  state is fully filled by electrons in the neutral state, while the  $e$  states remain empty and are able to accept up to four excess electrons. These  $e$  states serve as the localization center for electrons, enabling the formation of hyper-reduced states.

In higher amount of vacancies, the structural distortion remains to be similar to that for the isolated vacancy detailed in the above. For example, the neighboring Mo ions are displaced inward by only 0.03 Å even at the largest defect concentration in the present study (~25%). About the electronic structure, more defect levels are introduced in higher vacancy concentrations (see Figure S1 below; the  $a_1$ - and  $e$ -derived states are marked in the green and red circles, respectively). These levels are still contributed by Mo ions neighboring the vacancy sites (see partial charge densities below). Since the electronic structure is essentially identical regardless of the vacancy density, we expect that the hyper-reduced state will also play an important role in the catalytic reaction at high vacancy concentrations, but the detailed thermodynamic and kinetic parameters could be different.



**Figure S1.** The density of states (DOS) of vacant MoS<sub>2</sub> with 25% of the S-vacancy concentration in the upper S plane. The  $a_1$  related states are shown in green circle and  $e$  states in red circle. The corresponding partial charge density is also provided.

## Free energy calculation

The free energy  $G(n, q, N_{\text{H}_2})$  of the system in Figure 1b where  $N_{\text{H}_2}$  molecules of  $\text{H}_2$  evolved and the defect site is bonded to  $n$  hydrogen atoms ( $0 \leq n \leq 2$ ) and charged with  $q$ , can be written as follows:

$$G(n, q, N_{\text{H}_2}) - G(0,0,0) = G([n\text{H} - V_s]^q) + N_{\text{H}_2}\mu_{\text{H}_2} - G([V_s]^0) - (n + 2N_{\text{H}_2} - q)\mu_e - (n + 2N_{\text{H}_2})\mu_{\text{H}^+} \quad , \quad (1)$$

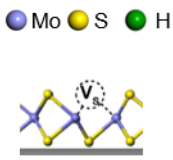
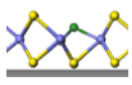
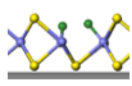
where  $G(0,0,0)$  is the reference state. In eq 1,  $(n + 2N_{\text{H}_2} - q)$  and  $(n + 2N_{\text{H}_2})$  are differences in the electron and proton number with respect to the reference state, respectively. The eq 1 can be simplified by applying following equations:

$$\mu_e^{\text{SHE}} + \mu_{\text{H}^+} = \mu_{\text{H}_2}, \quad (2)$$

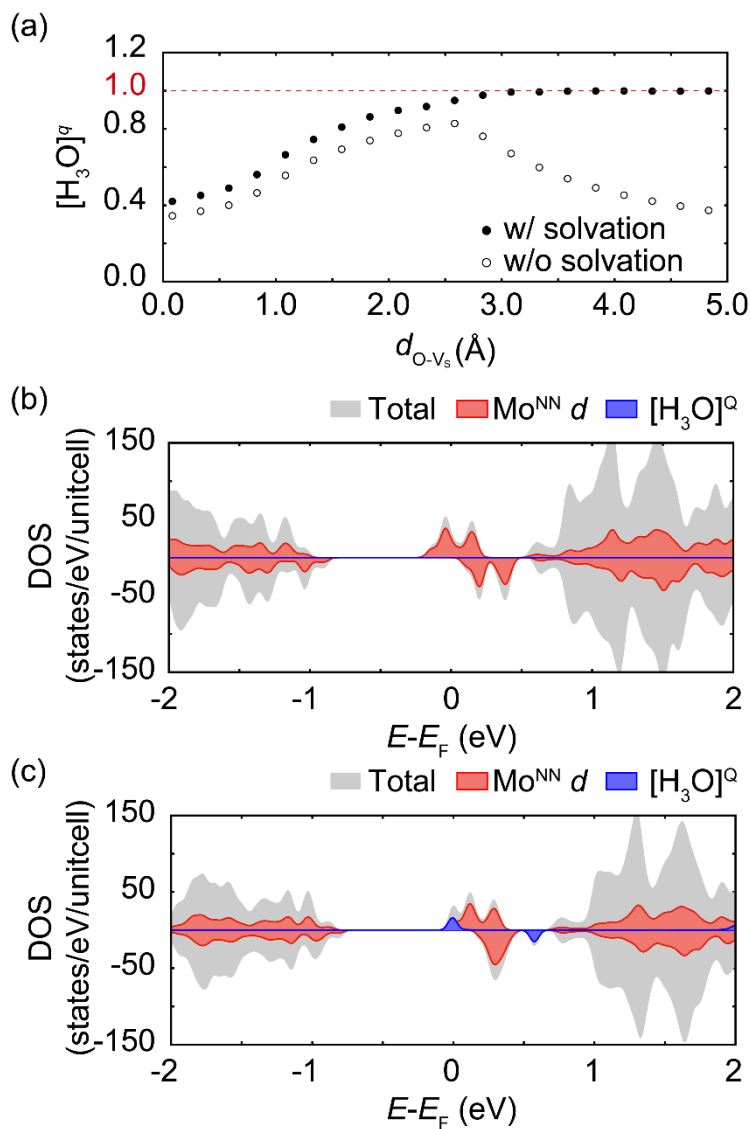
$$\mu_e = \mu_e^{\text{SHE}} + eU, \quad (3)$$

where  $e$  is the electron charge ( $< 0$ ). The eq 2 corresponds to the SHE condition, assuming that the proton source of the reaction is only hydronium ions. The eq 3 represents the chemical potential of electron under the overpotential  $U$  applied to the cathode. In the SHE condition,  $\mu_e^{\text{SHE}} = -4.44$  eV. By removing  $\mu_{\text{H}_2}$  and  $\mu_{\text{H}^+}$  in eq. 1 using eqs 2 and 3, one obtains the following equation:

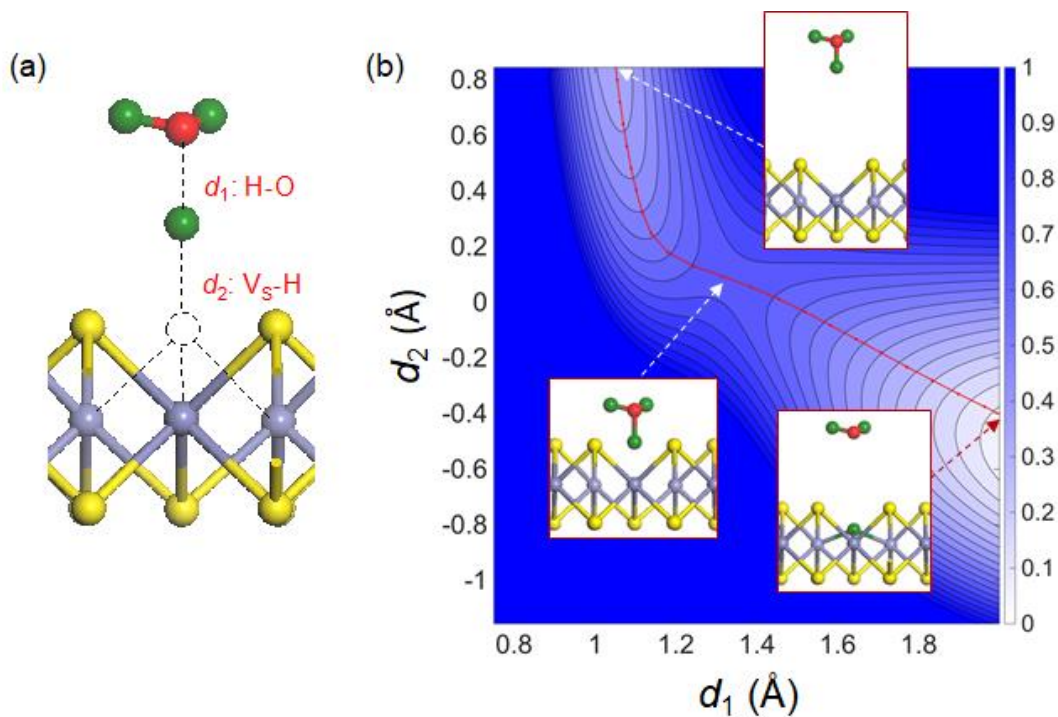
$$G(n, q, N_{\text{H}_2}) - G(0,0,0) = G([n\text{H} - V_s]^q) - G([V_s]^0) + q\mu_e^{\text{SHE}} - \frac{n}{2}\mu_{\text{H}_2} - eU(n + 2N_{\text{H}_2} - q) \quad (4).$$

Step		Reaction formula	$\Delta G$ (eV)
$[\mathcal{V}_S]^q \rightarrow [\mathcal{V}_S]^{q-1}$		$[\mathcal{V}_S]^{2+} + e^- \rightarrow [\mathcal{V}_S]^+$	$-1.356 + eU$
		$[\mathcal{V}_S]^+ + e^- \rightarrow [\mathcal{V}_S]$	$-1.219 + eU$
		$[\mathcal{V}_S] + e^- \rightarrow [\mathcal{V}_S]^-$	$-0.049 + eU$
		$[\mathcal{V}_S]^- + e^- \rightarrow [\mathcal{V}_S]^{2-}$	$0.286 + eU$
		$[\mathcal{V}_S]^{2-} + e^- \rightarrow [\mathcal{V}_S]^{3-}$	$0.622 + eU$
		$[\mathcal{V}_S]^{3-} + e^- \rightarrow [\mathcal{V}_S]^{4-}$	$0.780 + eU$
$[\text{H-}\mathcal{V}_S]^q \rightarrow [\text{H-}\mathcal{V}_S]^{q-1}$		$[\text{H-}\mathcal{V}_S]^+ + e^- \rightarrow [\text{H-}\mathcal{V}_S]^0$	$-0.511 + eU$
		$[\text{H-}\mathcal{V}_S]^0 + e^- \rightarrow [\text{H-}\mathcal{V}_S]^-$	$-0.185 + eU$
		$[\text{H-}\mathcal{V}_S]^- + e^- \rightarrow [\text{H-}\mathcal{V}_S]^{2-}$	$0.381 + eU$
		$[\text{H-}\mathcal{V}_S]^{2-} + e^- \rightarrow [\text{H-}\mathcal{V}_S]^{3-}$	$0.616 + eU$
$[2\text{H-}\mathcal{V}_S]^q \rightarrow [2\text{H-}\mathcal{V}_S]^{q-1}$		$[2\text{H-}\mathcal{V}_S]^0 + e^- \rightarrow [2\text{H-}\mathcal{V}_S]^-$	$-0.058 + eU$
		$[2\text{H-}\mathcal{V}_S]^- + e^- \rightarrow [2\text{H-}\mathcal{V}_S]^{2-}$	$0.388 + eU$

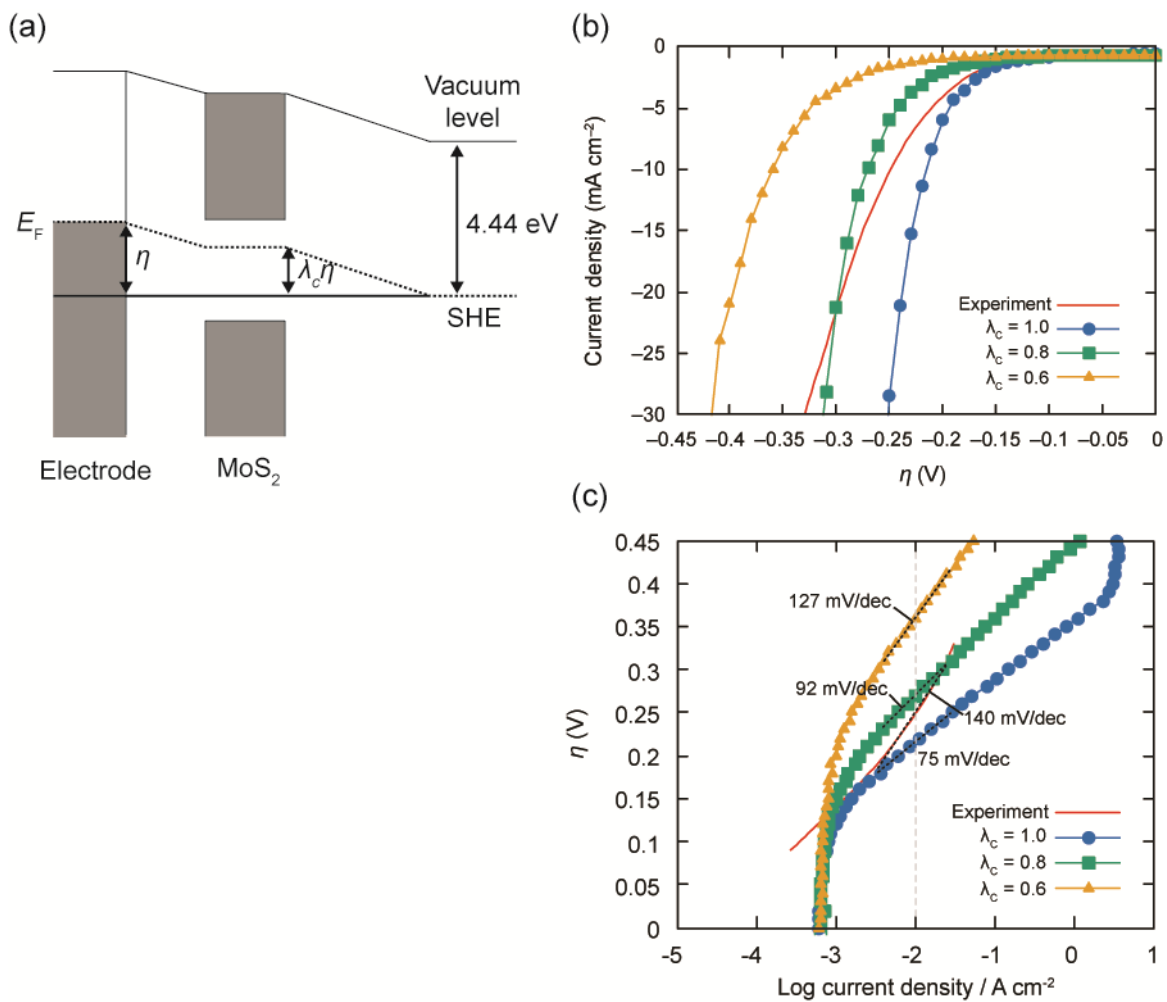
**Table S1.** The change in the Gibbs free energy during the charging step.



**Figure S2.** (a) Calculated charge  $q$  of a hydronium ion,  $[H_3O]^q$ , depending on the distance between  $[H_3O]^q$  and  $V_S$  with or without solvation effects. The model system includes MoS<sub>2</sub> substrate with one  $V_S$  and one electron is removed from the neutral charge state of the whole system. It is seen that +1 charged state of  $[H_3O]^q$  in water is properly described only when the solvation effects are included. (b), (c) Density of states (DOS) when the distance between  $[H_3O]^q$  and  $V_S$  is 5 Å (b) with or (c) without solvation effects.



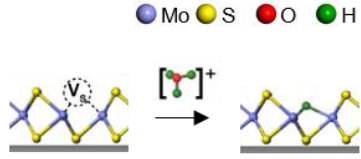
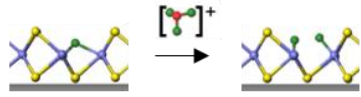
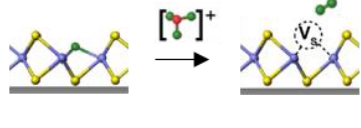
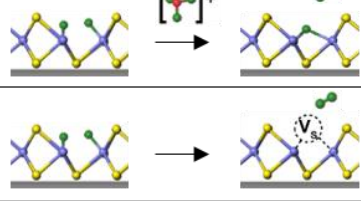
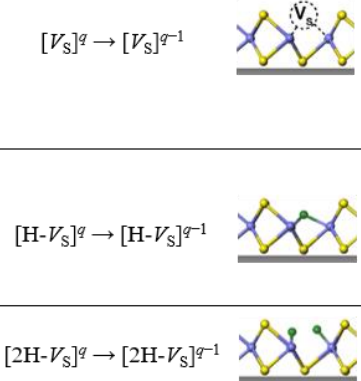
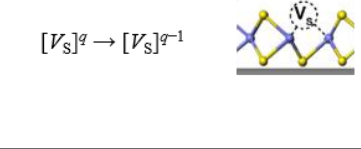
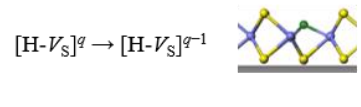
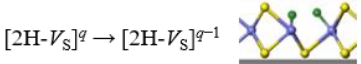
**Figure S3.** Since the NEB method in conjunction with solvation effects converge very slowly, we obtain the potential energy surface by the constrained minimization and identify the minimum energy path using the string method. (a) Constrained variables  $d_1$  and  $d_2$  in a Volmer step. (b) Potential energy surface of  $[V_S]^- + H_3O^+ \rightarrow [H-V_S]^0 + H_2O$  with respect to constrained variables of  $d_1$  and  $d_2$ . The minimum energy path identified by the string method is noted as red lines. The initial, final, and transition states are shown as insets.



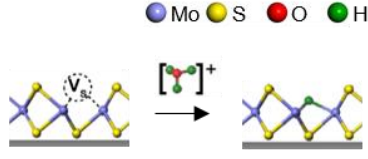
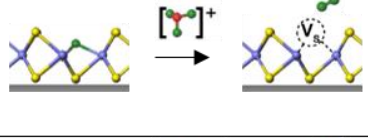

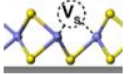
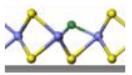
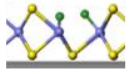
**Figure S4.** (a) Schematic diagram showing that the overpotential applied on MoS<sub>2</sub> ( $\eta_{\text{MoS}_2}$ ) is lower than for electrode ( $\eta$ ) due to the bad contact. As a result,  $\eta_{\text{MoS}_2} = \lambda_c \cdot \eta$  ( $\lambda_c \leq 1$ ) (b, c) Influence of  $\lambda_c$  on the (b) polarization curve, and (c) Tafel plot.

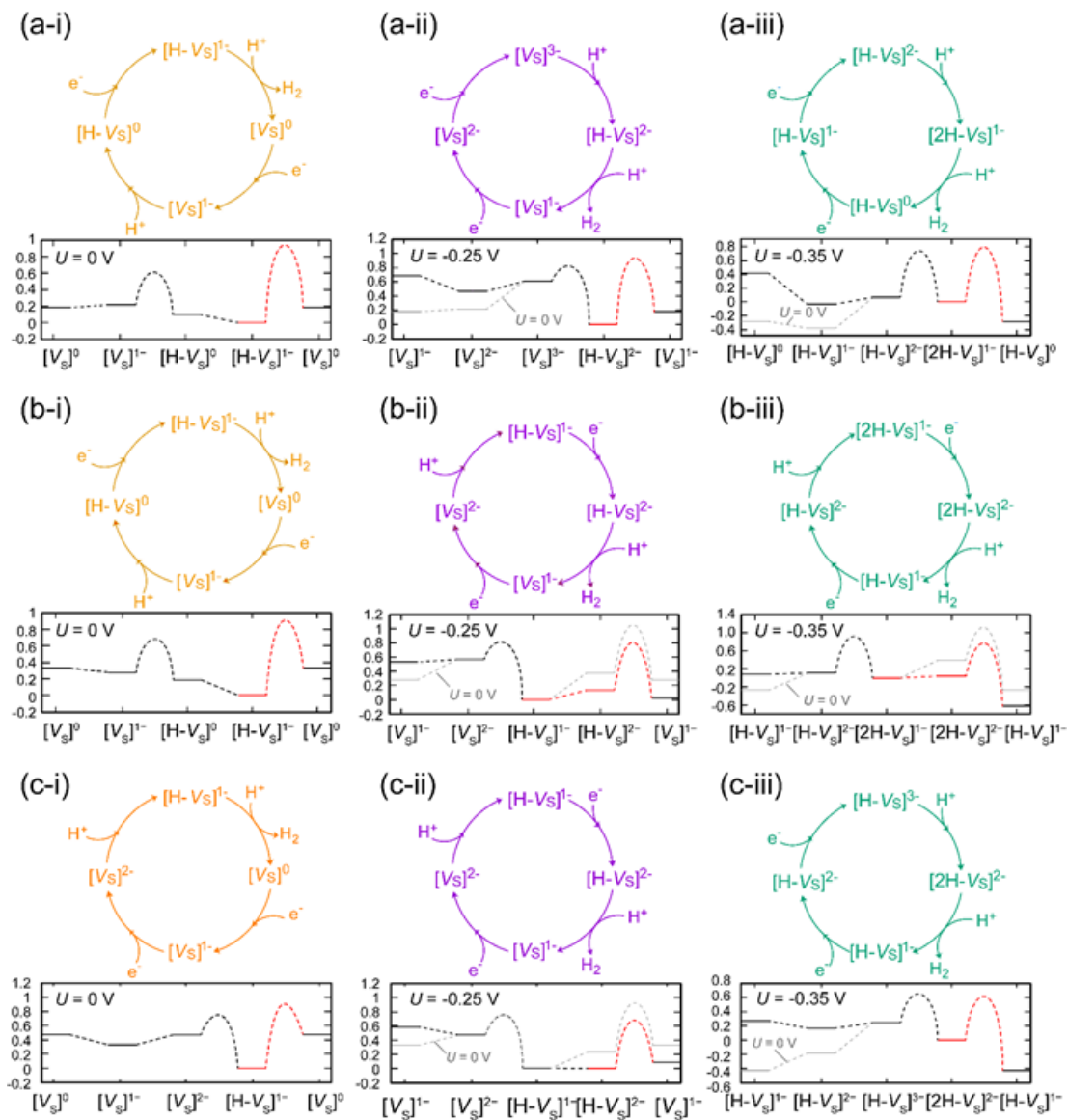


**Table S2.**  $\Delta G$  and  $\Delta G^\ddagger$  for elementary steps under the compressive strain of  $-1.35\%$ .

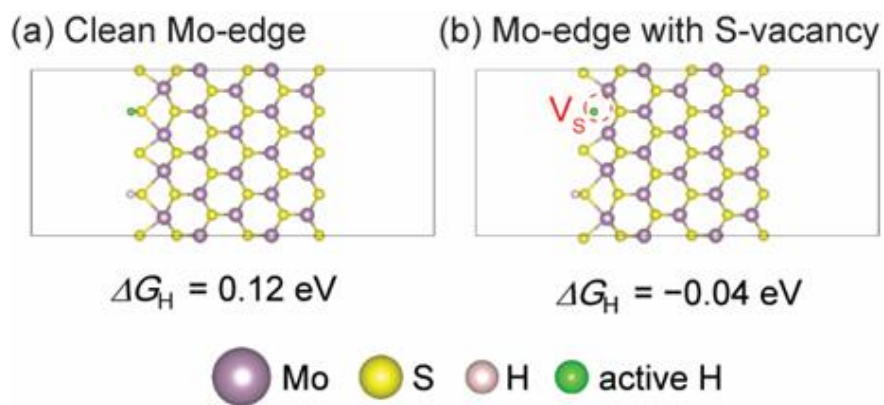
Step		Reaction formula	$\Delta G$ (eV)	$\Delta G^\ddagger$ (eV)		
Volmer		$[V_s] + H^+ \rightarrow [H-V_s]^+$	0.337	0.565		
		$[V_s]^- + H^+ \rightarrow [H-V_s]^0$	-0.141	0.391		
		$[V_s]^{2-} + H^+ \rightarrow [H-V_s]^-$	-0.649	0.191		
		$[V_s]^{3-} + H^+ \rightarrow [H-V_s]^{2-}$	-0.947	0.086		
		$[V_s]^{4-} + H^+ \rightarrow [H-V_s]^{3-}$	-1.106	0.000		
		$[H-V_s] + H^+ \rightarrow [2H-V_s]$	0.420	1.028		
		$[H-V_s]^- + H^+ \rightarrow [2H-V_s]^-$	-0.021	0.694		
		$[H-V_s]^{2-} + H^+ \rightarrow [2H-V_s]^{2-}$	-0.249	0.408		
		Heyrovsky		$[H-V_s]^+ + H^+ \rightarrow [V_s]^{2+} + H_2$	2.163	1.298
				$[H-V_s]^0 + H^+ \rightarrow [V_s]^+ + H_2$	1.195	0.921
$[H-V_s]^- + H^+ \rightarrow [V_s]^0 + H_2$	0.036			0.971		
$[H-V_s]^{2-} + H^+ \rightarrow [V_s]^- + H_2$	-0.359			1.120		
$[H-V_s]^{3-} + H^+ \rightarrow [V_s]^{2-} + H_2$	-0.631			0.471		
Tafel		$[2H-V_s]^0 + H^+ \rightarrow [H-V_s]^+ + H_2$	-0.047	0.916		
		$[2H-V_s]^- + H^+ \rightarrow [H-V_s]^0 + H_2$	-0.479	0.811		
		$[2H-V_s]^{2-} + H^+ \rightarrow [H-V_s]^- + H_2$	-1.031	0.727		
		Charging		$[2H-V_s]^0 \rightarrow [V_s]^0 + H_2$	-0.384	1.318
				$[2H-V_s]^- \rightarrow [V_s]^- + H_2$	-0.338	1.256
$[2H-V_s]^{2-} \rightarrow [V_s]^{2-} + H_2$	-0.382			1.299		
$[V_s]^{2+} + e^- \rightarrow [V_s]^+$	$-1.330 + eU$					
$[V_s]^+ + e^- \rightarrow [V_s]$	$-1.170 + eU$					
$[V_s] + e^- \rightarrow [V_s]^-$	$0.116 + eU$					
Charging		$[V_s]^- + e^- \rightarrow [V_s]^{2-}$	$0.497 + eU$			
		$[V_s]^{2-} + e^- \rightarrow [V_s]^{3-}$	$0.809 + eU$			
		$[V_s]^{3-} + e^- \rightarrow [V_s]^{4-}$	$0.928 + eU$			
		Charging		$[H-V_s]^+ + e^- \rightarrow [H-V_s]^0$	$-0.362 + eU$	
$[H-V_s]^0 + e^- \rightarrow [H-V_s]^-$	$-0.011 + eU$					
$[H-V_s]^- + e^- \rightarrow [H-V_s]^{2-}$	$0.511 + eU$					
$[H-V_s]^{2-} + e^- \rightarrow [H-V_s]^{3-}$	$0.769 + eU$					
Charging		$[2H-V_s]^0 + e^- \rightarrow [2H-V_s]^-$	$0.070 + eU$			
		$[2H-V_s]^- + e^- \rightarrow [2H-V_s]^{2-}$	$0.541 + eU$			

**Table S3.**  $\Delta G$  and  $\Delta G^\ddagger$  for elementary steps under the tensile strain of 1.35%.

Step	Reaction formula	$\Delta G$ (eV)	$\Delta G^\ddagger$ (eV)		
Volmer		$[V_s] + H^+ \rightarrow [H-V_s]^+$	0.417	0.609	
	$[V_s]^- + H^+ \rightarrow [H-V_s]^0$	0.014	0.485		
	$[V_s]^{2-} + H^+ \rightarrow [H-V_s]^-$	-0.370	0.314		
	$[V_s]^{3-} + H^+ \rightarrow [H-V_s]^{2-}$	-0.500	0.202		
	$[V_s]^{4-} + H^+ \rightarrow [H-V_s]^{3-}$	-0.632	0.000		
	$[H-V_s]^- + H^+ \rightarrow [2H-V_s]^0$	0.408	1.024		
	$[H-V_s]^{2-} + H^+ \rightarrow [2H-V_s]^-$	-0.007	0.644		
	$[H-V_s]^{3-} + H^+ \rightarrow [2H-V_s]^{2-}$	-0.161	0.415		
	Heyrovsky		$[H-V_s]^+ + H^+ \rightarrow [V_s]^{2+} + H_2$	1.846	1.217
		$[H-V_s]^0 + H^+ \rightarrow [V_s]^+ + H_2$	1.290	1.027	
$[H-V_s]^- + H^+ \rightarrow [V_s]^0 + H_2$		0.619	0.902		
$[H-V_s]^{2-} + H^+ \rightarrow [V_s]^- + H_2$		0.279	0.695		
$[H-V_s]^{3-} + H^+ \rightarrow [V_s]^{2-} + H_2$		0.020	0.573		
$[2H-V_s]^0 + H^+ \rightarrow [H-V_s]^+ + H_2$		0.628	0.904		
$[2H-V_s]^- + H^+ \rightarrow [H-V_s]^0 + H_2$		0.300	0.822		
$[2H-V_s]^{2-} + H^+ \rightarrow [H-V_s]^- + H_2$		-0.189	0.586		
Tafel			$[2H-V_s]^0 \rightarrow [V_s]^0 + H_2$	0.211	1.242
		$[2H-V_s]^- \rightarrow [V_s]^- + H_2$	0.286	1.206	
	$[2H-V_s]^{2-} \rightarrow [V_s]^{2-} + H_2$	0.181	1.298		
Charging	$[V_s]^q \rightarrow [V_s]^{q-1}$ 	$[V_s]^{2+} + e^- \rightarrow [V_s]^+$	-1.199+eU		
		$[V_s]^+ + e^- \rightarrow [V_s]$	-1.064+eU		
		$[V_s] + e^- \rightarrow [V_s]^-$	-0.240+eU		
		$[V_s]^- + e^- \rightarrow [V_s]^{2-}$	-0.009+eU		
		$[V_s]^{2-} + e^- \rightarrow [V_s]^{3-}$	0.230+eU		
		$[V_s]^{3-} + e^- \rightarrow [V_s]^{4-}$	0.382+eU		
	$[H-V_s]^q \rightarrow [H-V_s]^{q-1}$ 	$[H-V_s]^+ + e^- \rightarrow [H-V_s]^0$	-0.643+eU		
		$[H-V_s]^0 + e^- \rightarrow [H-V_s]^-$	-0.393+eU		
		$[H-V_s]^- + e^- \rightarrow [H-V_s]^{2-}$	0.100+eU		
		$[H-V_s]^{2-} + e^- \rightarrow [H-V_s]^{3-}$	0.250+eU		
	$[2H-V_s]^q \rightarrow [2H-V_s]^{q-1}$ 	$[2H-V_s]^0 + e^- \rightarrow [2H-V_s]^-$	-0.315+eU		
		$[2H-V_s]^- + e^- \rightarrow [2H-V_s]^{2-}$	0.096+eU		



**Figure S5.** Major paths on different conditions of strain ((a) compressive, (b) no strain, and (c) tensile strain) when applied potential is i) 0, ii)  $-0.25$  V, and iii)  $-0.35$  V, respectively. The red line represents the rate-limiting step and gray lines represent the corresponding path when the potential is 0.



**Figure S6.** Mo-edge configurations (a) with and (b) without S-vacancy. We take into account the hydrogen coverage of 0.25 in one side of edges.

**Table S4.** The dependence of the hydrogen binding energy ( $\Delta G_H$ ) of  $V_S$  on the number of MoS<sub>2</sub> layers.

Number of layers	$\Delta G_H$ (eV)
1	0.02
2	0.03
3	0.03
4	0.03